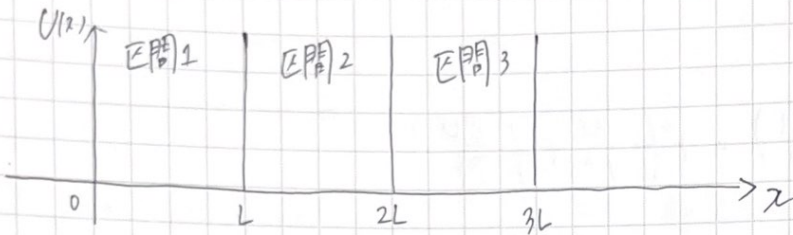


H=9 東工院試 午後

11) 
$$U(x) = a \sum_{n=-\infty}^{\infty} \delta(x - nL), \quad a > 0$$

$$\psi(x + 3L) = \psi(x)$$



$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + U(x)$$

$\hat{\zeta} =$  右シフト演算子

$$\hat{\zeta} \psi(x) = \psi(x - L)$$

$\hat{p} =$  反転演算子

$$\hat{p} \psi(x) = \psi(-x)$$

(1)

(3) 
$$\hat{\zeta}^{-1} \psi(x) = \psi(x + L)$$

(1) 恒等演算子

$$\hat{\zeta} \psi = \eta \psi$$

$$\hat{\zeta}^3 \psi = \eta^3 \psi = \psi$$

$$\eta^3 = 1$$

(1) 
$$\eta = 1, \omega, \omega^2$$

(2) 
$$\hat{p} \psi = \eta' \psi$$

$$\hat{p}^2 \psi = \eta'^2 \psi = \psi$$

$$\eta'^2 = 1$$

$$\eta' = 1, -1$$

(1) 交換可能

(2) 
$$\hat{p}^{-1} \hat{\zeta} \hat{p} \psi(x)$$

$$= \hat{p}^{-1} \hat{\zeta} \psi(-x)$$

~~$$= \hat{p}^{-1} \psi(-x - L)$$~~

~~$$= \psi(-x - L)$$~~

~~$$= \hat{\zeta}^{-1} \psi(x)$$~~

$$= \hat{p}^{-1} \psi(-(x - L))$$

$$= \hat{p}^{-1} \psi(-x + L)$$

$$= \psi(x - L)$$

$$= \hat{\zeta} \psi(x)$$

(3) 
$$\hat{p}^{-1} \hat{\zeta} \hat{p} = \hat{\zeta} \quad \forall \psi$$

$$\hat{\zeta} \hat{p} = \hat{p} \hat{\zeta} \quad \text{成り立つ}$$

$$\hat{\zeta} \psi = \eta \psi$$

$$\hat{p} \hat{\zeta} \psi = \eta \hat{p} \psi$$

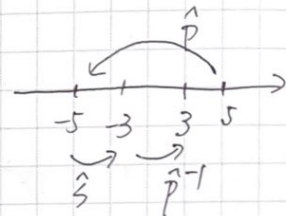
$$\hat{\zeta} \hat{p} \psi = \eta \hat{p} \psi$$

$\psi$  が  $\eta$  の固有関数と仮定

$$\therefore \hat{p} \psi \text{ の固有値} = 1, \omega, \omega^2$$

図は

$$x = 5, L = 2$$



$$5 \rightarrow 5 - 2$$

$$x \rightarrow x - L$$

$$\therefore \hat{\zeta} \psi(x) = \psi(x - L)$$

波動関数の係数  $a \rightarrow \infty$

$$\psi_1(x) = \begin{cases} \sqrt{\frac{2}{L}} \sin\left(\frac{\pi}{L}x\right) & (0 \leq x \leq L) \\ 0 & (L \leq x \leq 3L) \end{cases}$$

$$\psi_2(x) = \hat{\omega} \psi_1(x)$$

$$\psi_3(x) = \hat{\omega}^2 \psi_1(x)$$

$\omega$  が解根, 3つのエネルギー固有状態の波動関数

$$\Phi_1(x) = \frac{1}{\sqrt{3}} (\psi_1(x) + \psi_2(x) + \psi_3(x)) \Rightarrow E_1$$

$$\Phi_2(x) = \frac{1}{\sqrt{3}} (\psi_1(x) + \omega \psi_2(x) + \omega^2 \psi_3(x))$$

$$\Phi_3(x) = \frac{1}{\sqrt{3}} (\psi_1(x) + \omega^2 \psi_2(x) + \omega \psi_3(x))$$

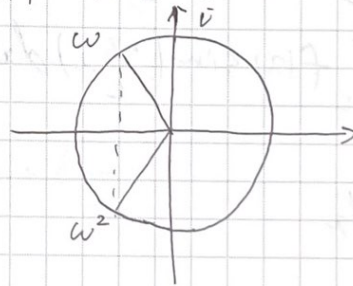
(2)  $\psi_1 + \psi_2 + \psi_3$  の基底状態の線形結合  $\Psi$

$E_1$  の 1 番低い

$$\omega = e^{i\frac{2}{3}\pi} \quad \omega^2 = e^{i\frac{4}{3}\pi}$$

$$\omega^* = \omega^2$$

$$\therefore \Phi_3 = \Phi_2^*$$



$$\hat{H}\Phi_2 = E_2\Phi_2$$

$$\hat{H}\Phi_3 = E_3\Phi_3 = \hat{H}\Phi_2^* = E_3\Phi_2^*$$

$$\Phi_2^* \hat{H} \Phi_2 = \Phi_2^* E_2 \Phi_2$$

$$\langle \Phi_2 | \hat{H} | \Phi_2 \rangle = E_2 |\Phi_2|^2$$

$$\Phi_2 \hat{H} \Phi_2^* = \Phi_2 E_3 \Phi_2^*$$

$$\langle \Phi_2 | \hat{H} | \Phi_2 \rangle^* = E_3 |\Phi_2|^2$$

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + U(x) \quad (\because U(x) = a \sum_{n=-\infty}^{\infty} \delta(x-nL), a > 0)$$

実部のみを  $\hat{H}$  は  $\hat{H}$  である。

$$\langle \Phi_2 | \hat{H} | \Phi_2 \rangle^* = \langle \Phi_2 | \hat{H} | \Phi_2 \rangle \quad (xL \equiv -L)$$

$$\therefore E_2 = E_3 \quad (A)$$